



IMPACT OF FILAMENTS ON THE RELATIVE DISPERSION OF SURFACE DRIFTERS IN THE BENGUELA UPWELLING REGION

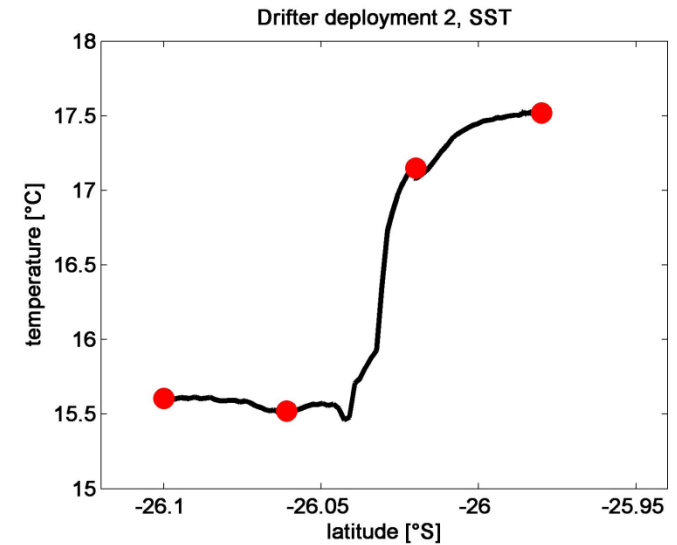
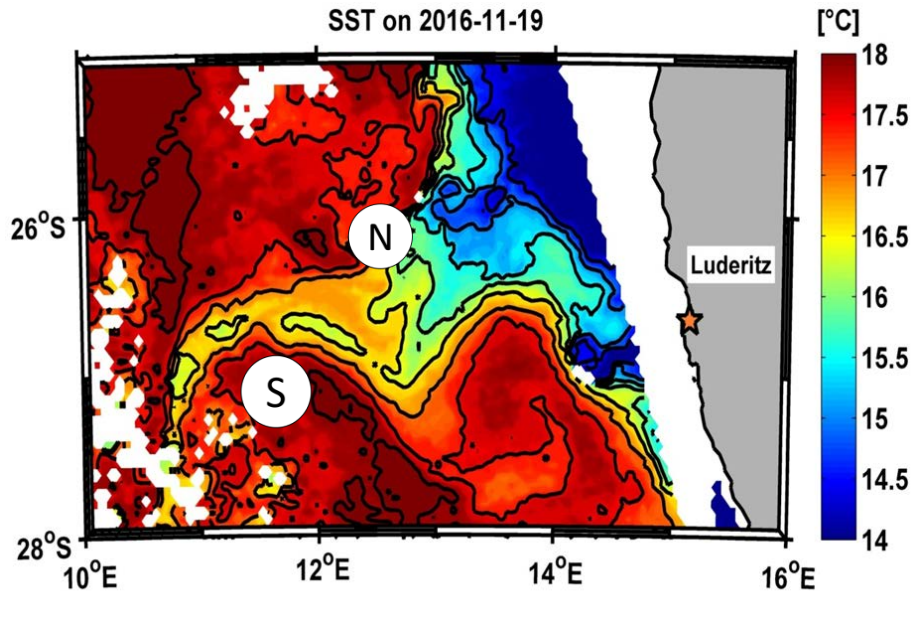
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Overview

- Surface drifter experiment in the Benguela upwelling region
- (Relative) dispersion of pairs of drifters:
 - fixed times analysis: mean squared pair separations
 - fixed scale analysis: Finite Size Lyapunov Exponents (FSLEs)
- Mean squared pair separations: Different dispersion regimes
- Anomalous diffusion
- Do Finite Size Lyapunov Exponents yield consistent results ?
- Analysis of drifter subsets: dependence on deployment location
- Do our data show local (Richardson) dynamics?
- Conclusion



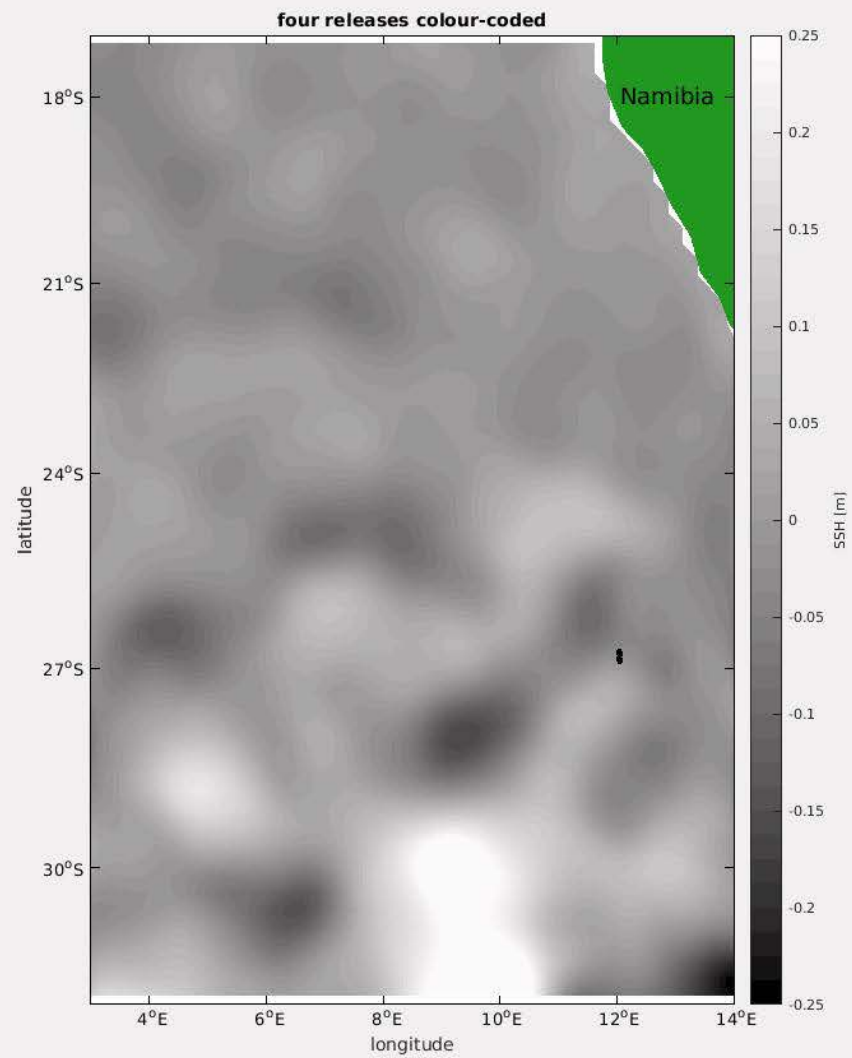
Upwelling water front off Namibia's coast with filament



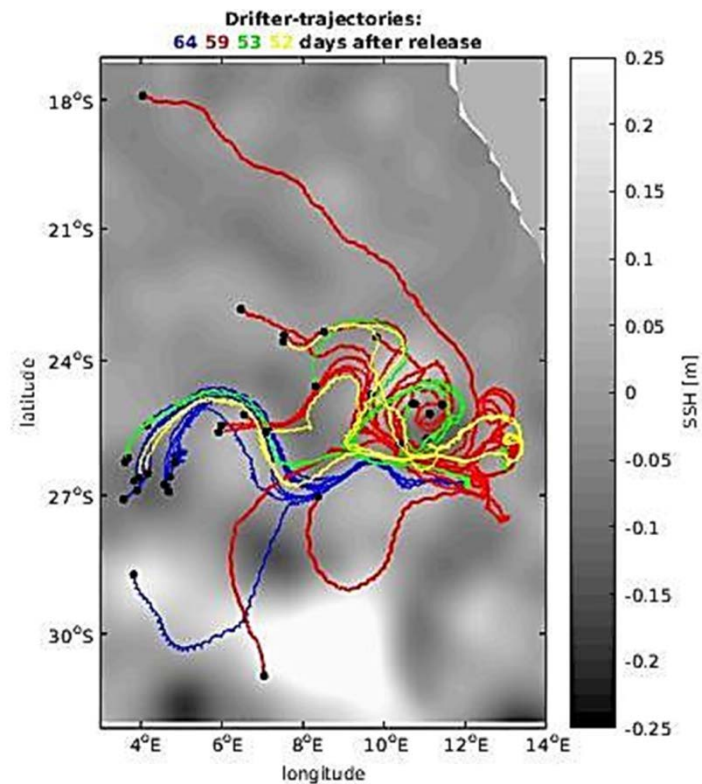
*SST along ship track.
Red dots: triplet-deployments*

36 surface drifters were released in 4 groups (12/12/6/6 drifters) of triplets, with initial distance of 100 m –500 m. Deployment locations: southern boundary (27°S, 12°E) (S), northern boundary (26°S, 13°E) (N)





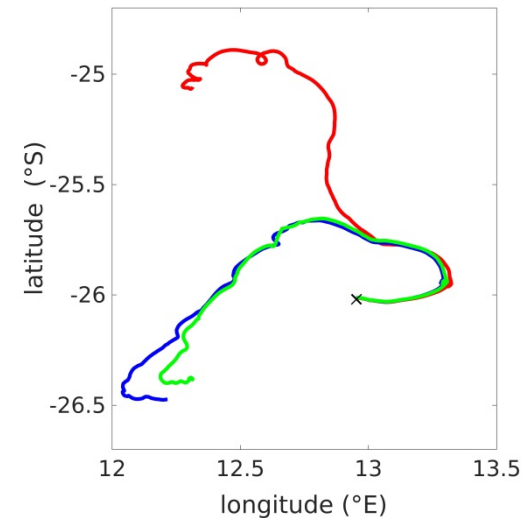
Drifter trajectories from cruise M132 in the area off Namibia's coast in november/december 2016



Lagrangian analysis

Pair separation statistics yield information about the dominant physical mechanisms acting at the different scales of motion (e.g. chaotic advection, turbulence, diffusion) in specific regions

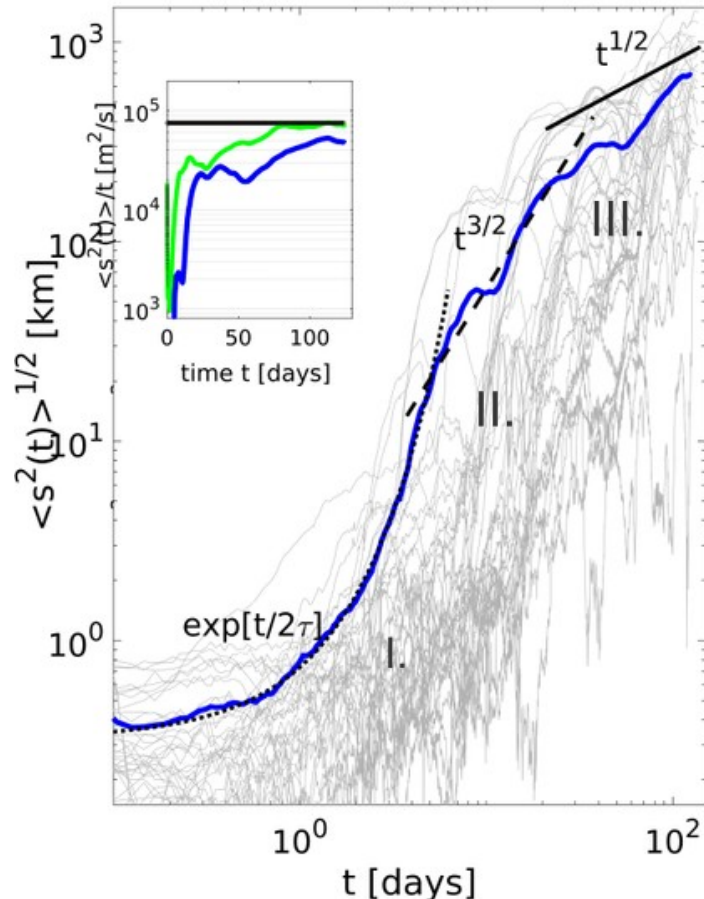
- fixed time measure:
Mean squared pair separation $\langle s^2(t) \rangle$
- Fixed scale measure:
Finite size Lyapunov exponents $\lambda(\delta) = \frac{1}{\langle \tau(\delta) \rangle} \ln(\alpha)$
 $\langle \tau(\delta) \rangle$ = averaged time for drifter-pairs, initially separated by distance δ , to reach distance $\alpha\delta$



What dispersion regimes can be detected (at what scales) ?
Do the two methods yield consistent results ?



Lagrangian root mean squared pair separation $\langle s^2(t) \rangle^{1/2}$



	pair separations
I. nonlocal	$s \sim e^{t/\tau}$, $\tau \cong 0.6$ days
II. local	$\langle s^2 \rangle = C\epsilon t^3$ $C\epsilon \cong 4.4 \times 10^{-9}$ W/kg
III. diffusive	$\langle s^2 \rangle = 2\kappa t$, $\kappa \cong 3.7 \times 10^4$ m ² /s

→ Nonlocal behavior for $s < \text{Rossby Radius } R_1 (\cong 30 \text{ km})$



- **Brownian motion** (Einstein, Ann. d. Phys., 1905)
- Deviation (diffusional anomalies)

$$\langle r^2(t) \rangle = 2\kappa t$$

$\alpha > 1$ (superdiffusion), $\kappa(t) \rightarrow t^\nu$

$\alpha < 1$ (subdiffusion), $\kappa(t) \rightarrow 0$

$$\langle r^2(t) \rangle \sim t^\alpha$$

experimentally observed in:

- $\alpha > 1$ floating tracers in ocean (Osborne et al., Tellus ' 89)
rotating fluid flows (Solomon.. Phys. Rev. Lett. ' 93),
sea ice (Gabrielsi, Badin, Kaleschke, J. o. Geoph. Res. ' 14)
- $\alpha < 1$ polymer network (Amblard et al., Phys. Rev. Lett. ' 96)
charge transport (Scher & Montroll Phys. Rev. B)

``Rules of thumb`` for physical mechanism:

- $\alpha > 1$: persistence of free paths (``Levy Flights``)
- $\alpha < 1$: geometric constraints (disorder or fractal)

→ strong diffusional anomaly

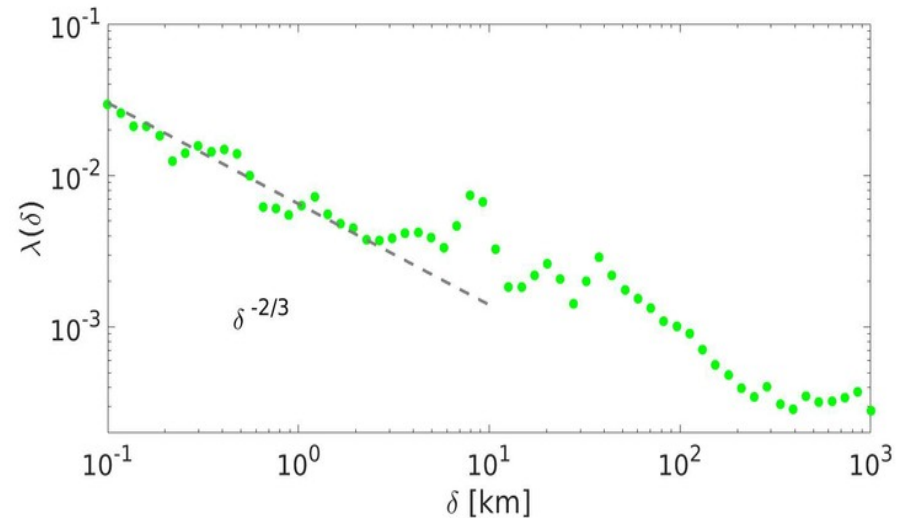
$$\langle r^2(t) \rangle \sim \ln^\alpha(t)$$

(Dräger & Klafter
Phys. Rev. Lett. ' 00)



Finite Size Lyapunov Exponents

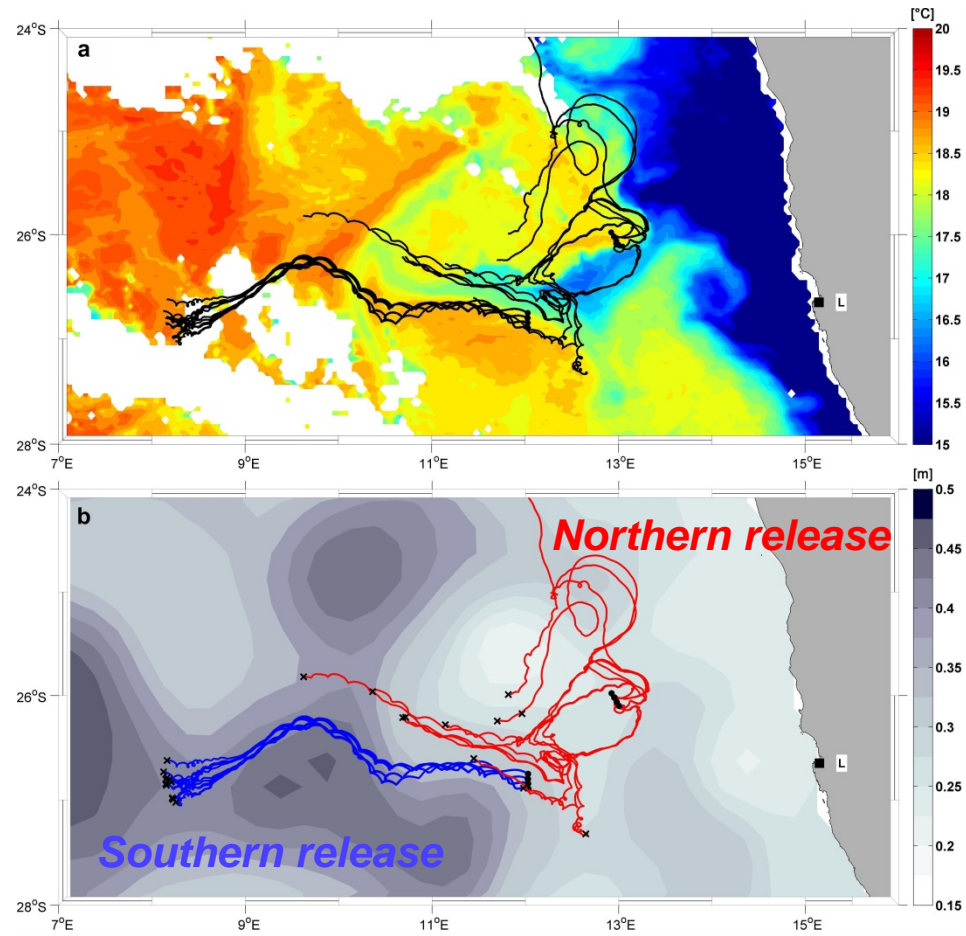
	pair separations	FSLE (expected)
I.Nonlocal	$s \sim e^{t/\tau}$	$\lambda(\delta) \sim \text{const}$
II.Local	$\langle s^2 \rangle \sim t^3$	$\lambda(\delta) \sim \delta^{-2/3}$
III.Diffusive	$\langle s^2 \rangle \sim t$	$\lambda(\delta) \sim \delta^{-2}$



Controversial results for small scales:
 FSLEs $\lambda(\delta)$ point to local behavior, $\langle s^2(t) \rangle$ to nonlocal behavior

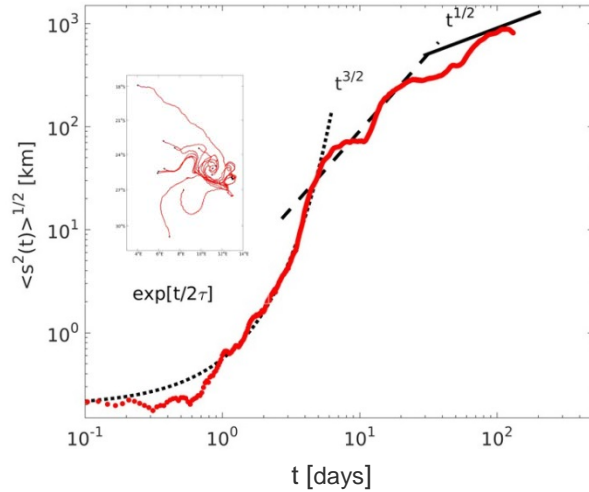


Different deployment scenarios:

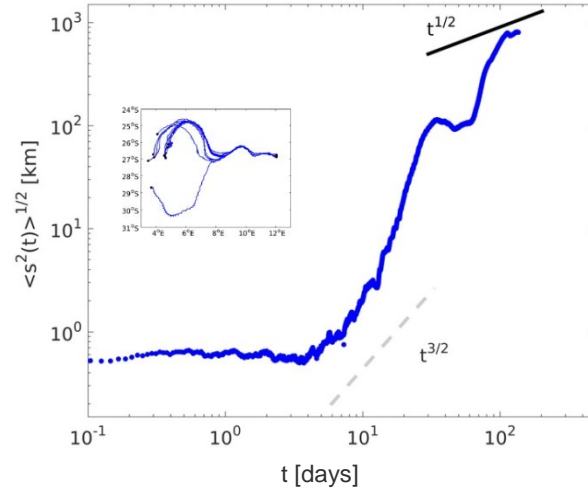


$$\langle s^2(t) \rangle^{1/2}$$

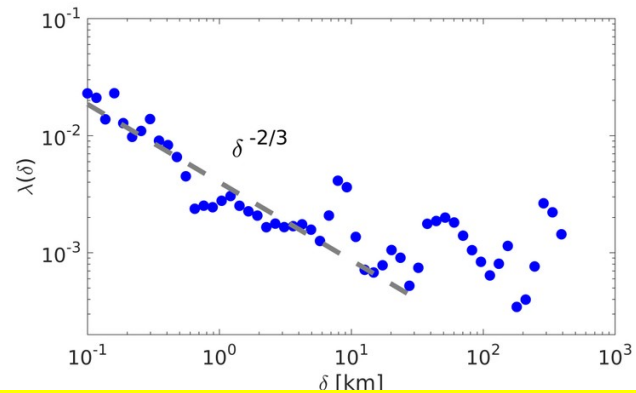
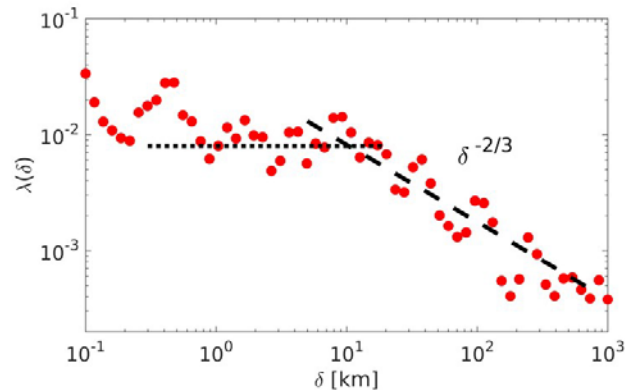
Northern release



Southern release



FSLE
 $\lambda(\delta)$



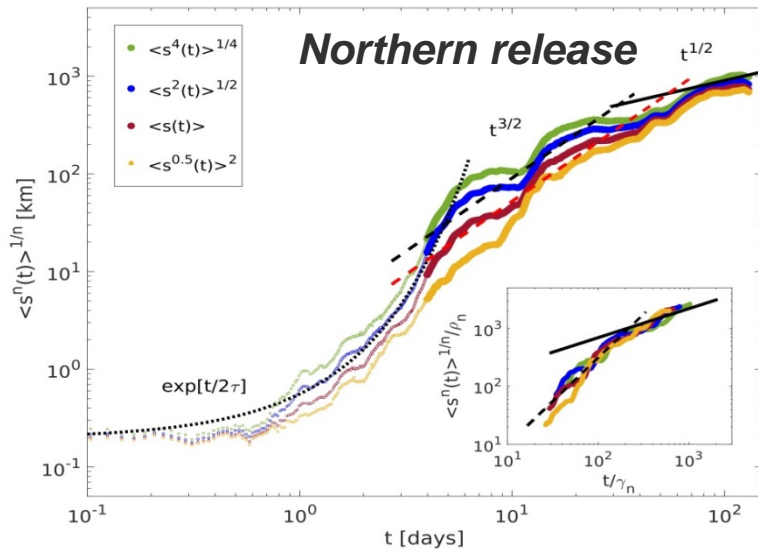
Consistent results for $\langle s^2(t) \rangle$ and $\lambda(\delta)$,
when the methods are applied to the two releases **separately**

Do our data show local (Richardson) dynamics? Below or above R_1 ? For both releases ?

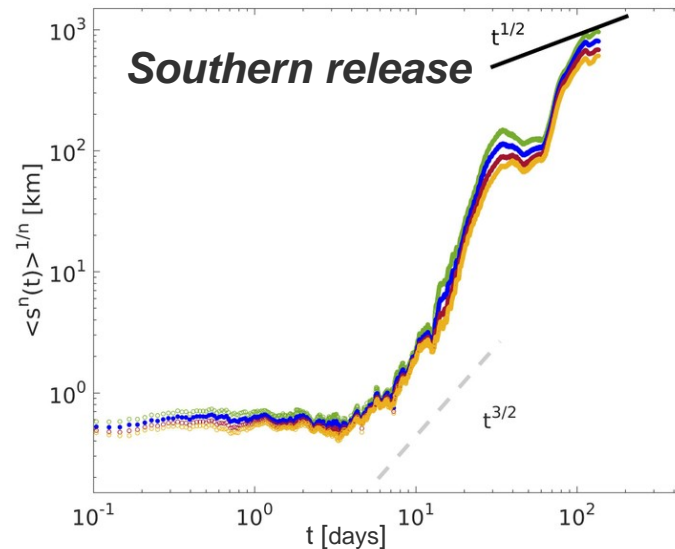
$\langle s^2(t) \rangle \sim \epsilon t^3$ for $s > R_1$ (**northern release** & whole ensemble)

$\lambda(\delta) \sim \delta^{-2/3}$ for $\delta \ll R_1$ (**southern release** & whole ensemble)

Hallmark of local transport: anomalous broad distribution: $P(s|t) = \frac{s}{3} \left(\frac{b}{t}\right)^3 \exp\left[\frac{-bs^\beta}{t}\right]$, $\beta = 2/3$
Would be confirmed by datacollapse of moments $\langle s^n \rangle$



YES



→ no intermittent
Richardson
dynamic

(Probably)
No



Conclusions

- The Lagrangian squared pair separations $\langle s^2 \rangle$ **averaged over all trajectories** show a nonlocal behavior for the first 4 days followed by a (local) Richardson regime (t^3 - law) and a diffusive regime $\langle s^2 \rangle \sim t$.
- In contrast the Finite Size Lyapunov-exponents (FSLE) $\lambda(\delta)$ are not constant for small δ , as expected for nonlocal transport, but decay as $\delta^{-2/3}$.
- The inconsistency originates in the strong influence of the location of deployment leading to a mixture of dynamics:
 - $\langle s^2(t) \rangle$ is dominated by the northern (inertial) release due to scale separation
 - $\lambda(\delta)$ reveals the small scale behavior **within** the filament (southern release)
- While the anomalous PDF (northern release) supports local transport for $s > R_1$ the narrow PDF (southern release) for small s questions submesoscale local dispersion.

J. Dräger, K. Jochumsen, A. Griesel and G. Badin submitted to Journal of Physical Oceanography



Thank you!

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The moments $\langle s^n(t) \rangle^{1/n}$ can be derived analytically straightforwardly from the theoretical forecast for $P(s|t)$ for the local (5)) and the diffusive regime (7), leading to

$$\langle s^n(t) \rangle^{1/n} = \begin{cases} i_n(t/b)^{3/2} & \text{for } t < t_x(n) \\ k_n(2\kappa t)^{1/2} & \text{for } t > t_x(n). \end{cases} \quad (8)$$

Here $i_n = (0.5 \Gamma((\frac{n+2}{2}) 3))^{1/n}$ and $k_n = \Gamma(\frac{n+2}{2})^{1/n}$ where $\Gamma(x)$ is the Gamma-function and with b defined in Eq. (5).

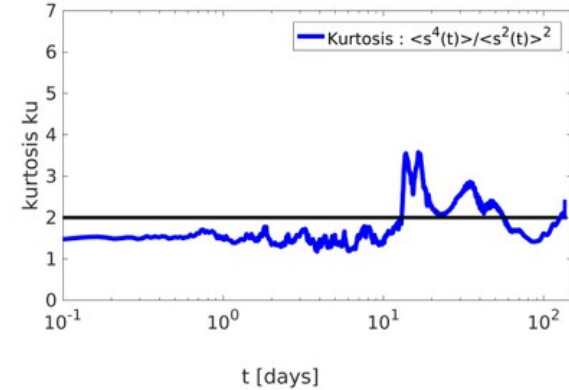
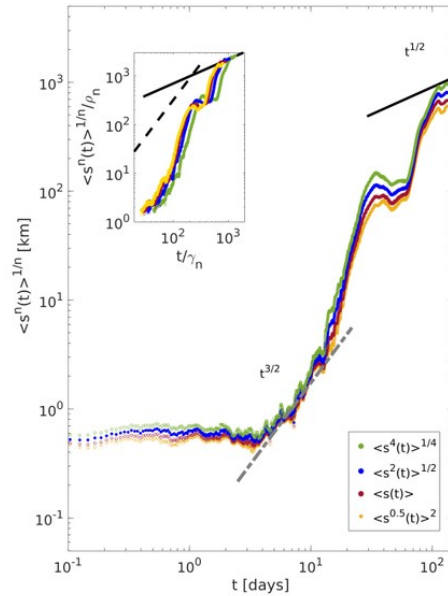
As the transition time $t_x(n)$ has to fulfill $i_n(t_x(n)/b)^{3/2} = k_n(2\kappa t_x(n))^{1/2}$, it is given by

$$t_x(n) = (2\kappa b^3)^{1/2} \gamma_n \quad (9)$$

with $\gamma_n = k_n/i_n$. Equations (8) and (9) indicate that the transition time $t_x(n)$ from the local to the diffusive dispersion regime differs for the different moments $\langle s^n(t) \rangle^{1/n}$.



Local submesoscale features within the filament ?

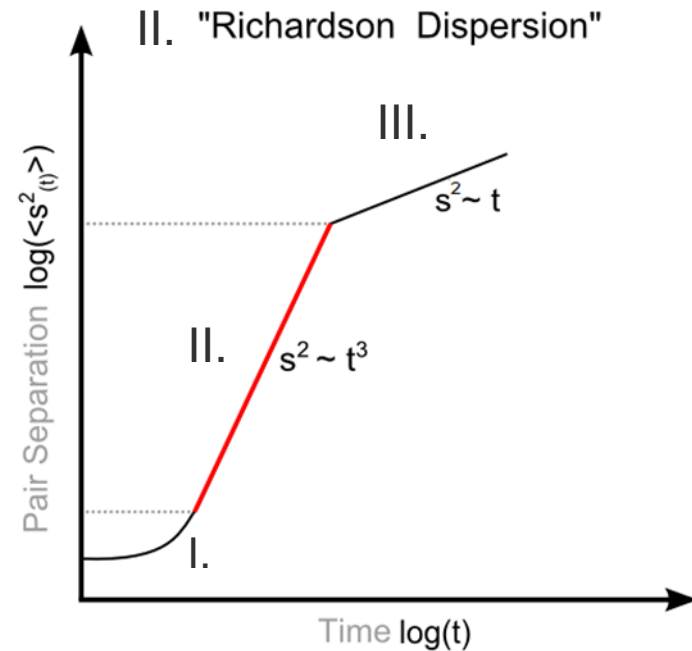


→ no intermittent Richardson dynamic



Mean squared pair separation: theoretical prediction

	fixed time measurement	fixed scale Measurement
I. nonlocal: s < eddies	$s \sim e^{t/\tau}$ $\tau = \text{unfolding rate}$	$\lambda(\delta) = \text{const}$
II. local: (Richardson) s \cong eddies	$\langle s^2 \rangle = \epsilon t^3$ $\epsilon = \text{energy transfer rate}$ $P(s t) \sim e^{-bs^{2/3}/t}$	$\lambda(\delta) \sim \delta^{-2/3}$
III. diffusive s > eddies	$\langle s^2 \rangle = 2\kappa t$ $\kappa = \text{relative diffusivity}$ $P(s t) \sim e^{-cs^2/t}$	$\lambda(\delta) \sim \delta^{-2}$



$\langle s^2(t) \rangle$: averaged squared separation of drifter-pairs, $\lambda(\delta)$: FSLE



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